# PREDICTION OF EMPEROR-HAWAII SEAMOUNT LOCATIONS FROM A REVISED MODEL OF GLOBAL PLATE MOTION AND MANTLE FLOW: 1. Method



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# **Concept:**

### **Physical model of mantle**

Assume relationship between density, temperature, depth, and seismic velocity. Viscosity model from Geoid and heat flow.

# Implant plumes at arbitrary time & place

Advection of plume conduits by large-scale flow and vertical buoyant rising. Fix to base of mantle or freely move.

# **Compute trace of plume at base of lithosphere**

**Predict present velocity field and past flow** Global plate motions and past locations of plate boundaries used as constraints in flow computation.

Iterate plume emplacement to optimise fit between predictions and known hotspot geology, and hence tie plate-motion model to mantle reference frame.

**Radial Mantle Viscosity structure**  $\eta(z)$ Stress-strain relationship  $\dot{\varepsilon} = C_1 \sigma^n exp\left(-\frac{H}{RT}\right)$  $\implies \eta(z) = \eta_0 \exp\left(\frac{H(z)}{nR\overline{T}(z)}\right) \cdot \left(\langle \dot{\varepsilon}^2 \rangle(z)\right)^{\frac{1-n}{2n}}$ 

 $\eta_0$  (may be different for upper mantle, transition zone, lower mantle) to be determined by optimizing fit to various observables (geoid, heat flux, CMB excess ellipticity ...)

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H = activation enthalpy T = temperature
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Empirical relationship:  $H = gRT_m$ ,  $T_m$  = melting temperature

Here use g/n = 12 (Yamazaki and Karato, 2001)









#### South Pacific plate-motion circuit







Obtained from optimizing fit to geoid, with additional heat flow constraint (Steinberger and Calderwood, 2001)

### Thermal expansivity profile

Integrate  $\frac{d \ln \alpha}{d \ln \rho} = -\delta_T$  along isotherms.

Mantle density anomalies

Seismic velocity anomalies  $\delta v_s$  from tomography Assume thermal origin of  $\delta v_s$  and density anomalies  $\delta \rho$ :  $F_{s,th} := (\delta \rho / \rho) / (\delta v_s / v_s) = (\alpha / \rho) / (\partial \ln v_s / \partial T)_p$ Orange line: g=12; Red line: g=27; Purple line: g=42



Relate seismic velocity and temperature anomalies Include anelastic and anharmonic effects:

 $-(\partial \ln v_s/\partial T)_p = -(\partial \ln v_{s,0}/\partial T)_p + (Q^{-1}/\pi) \cdot (gT_m/T^2)$ Dashed line: Anharmonic part. In the upper mantle after Goes et al. (2004); In the lower mantle, compute  $\mu(T, z)$ by integrating  $\frac{d\mu}{dz}(T_0(z) \pm \Delta T(z)) = \frac{d\mu}{dz}(T_0(z))(1 \pm \alpha(z)\Delta T(z))$  (1) (Duffy and Anderson, 1989). Starting point is Relative plate motion chains for times older than chron 20 (43 Ma). For times <43 Ma both models follow a chain through continental Antarctica, and include motion from Cande et al. (2000).

68 Ma reconstructions showing implications of alternate platemotion chains (83-43 Ma). Black arrows are scaled to be double the convergence in New Zealand (68-43 Ma), or rifting in Antarctica (68-26 Ma), that is predicted through closure of the South Pacific plate-motion circuit. Model 2 is in much better accord with local observations.



#### **Additional Cretaceous intra-plate deformation?**

Upper mantle:  $\delta_T = 5.5$ ; formalism by Schmeling et al. (2003).

Lower mantle:  $\delta_T = \delta_{T0} \left(\frac{\rho_0}{\rho}\right)^b$ 

 $\delta_{T0} = 5.5 \text{ and } b = 1.4$ 

 $\alpha_0(T) = (35 + 9T/1000 \text{K}) \cdot 10^{-6} \text{K}^{-1}$  for zero pressure

 $\rho(z)$  from PREM,  $\rho_0(z)$  from evaluating PREM lower mantle parameters at z=0 and accounting for temperature difference.



 $\mu(T_0(z_0) \pm \Delta T(z_0)) = \mu(T_0(z_0)) \pm \Delta T(z_0) \cdot \frac{d\mu_0}{dT}$ 

 $\frac{d\mu_0}{dT}$ =27 MPa/K;  $\frac{d\mu}{dz}(T_0(z))$  and  $\mu(T_0(z_0))$  from PREM. Dotted lines: Anelastic part. Continuous: Total Orange lines: g=12; Red lines: g=27; Purple lines: g=42



Although NOT inculuded in this analysis, additional intra-plate rifting that may affect both plate motion chains is suggested by New Zealand geology. Rifting is likely to have continued in places until chron 33 (79-74 Ma) and may have locally continued until as late as 60 Ma. Forward models show that implied block rotations have approximately the right geometry to improve our fit to the oldest part of the Emperor chain. The much greater local rotation of the Campbell Plateau, as compared to the long and thin Lord Howe Rise and New Caledonia Basin, means that model 1 is most severely affected by Cretaceous rifting.

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Details of other sources used are given in:

Steinberger, B., Sutherland, R., O'Connell, R.J. 2004. Prediction of Emperor-Hawaii seamount locations from a revised model of global plate motion and mantle flow. Nature 430: 167-173.

# PREDICTION OF EMPEROR-HAWAII SEAMOUNT LOCATIONS FROM A REVISED MODEL OF GLOBAL PLATE MOTION AND MANTLE FLOW: 2. Results & Conclusions



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N-S mantle cross section at 155° W

and projection of predicted Hawaiian plume conduit

10.

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(red has source that moves freely,

purple assumes fixed source)

Step 2:

insert

vertical

conduit

(here:

at time

170 Ma)

plume

conduit

gets

distorted

 $\mathbf{V}$ 

hotspot

moves

# **Conclusions (Nature 430: 167-173)**

Hotspot motion predicted by model of mantle flow improves fit for times <43 Ma

Predicted hotspot motion improves fit to Emperor seamount paleolatitudes 49-80 Ma, but does not improve longitude misfit

## Computed flow field

Based on model smean by Becker and Boschi (2002)  $(\delta\rho/\rho)/(\delta v_s/v_s)$ =0.2 below 220 km

→ 2 cm/yr



Plate motions uncertain before 43 Ma: 'missing' boundary in S Pacific

Model 1, no intra-Antarctic boundary before 43 Ma, produces poor fit to Emperor seamounts, and predicts convergence in New Zealand - inconsistent with rift structure

Model 2, no Tonga-Kermadec boundary before 43 Ma, produces a better fit to Emperor seamounts, and acceptably predicts Antarctic rift structure

Hotspot motion AND 'missing' intra-plate deformation are required to explain global observations of seamounts related to deep-seated hotspots

